

Carbon nanotube bolometers

M. Tarasov^{a)}

Institute of Radio Engineering and Electronics, 125009 Moscow, Russia

J. Svensson

Department of Physics, Göteborg University, SE-41296 Göteborg, Sweden

L. Kuzmin

Department of Microtechnology and Nanoscience, Chalmers University of Technology, SE-41296 Göteborg, Sweden

E. E. B. Campbell

Department of Physics, Göteborg University, SE-41296 Göteborg, Sweden

(Received 23 January 2007; accepted 13 March 2007; published online 16 April 2007)

A cryogenic bolometer has been fabricated using a bundle of single-walled carbon nanotubes as absorber. A bolometric response was observed when the device was exposed to radiation at 110 GHz. The temperature response was 0.4 mV/K, with an intrinsic electrical responsivity at low frequency up to 10^9 V/W and noise equivalent power of 3×10^{-16} W/Hz^{1/2} at 4.2 K. The response is largest at input power levels of a few femtowatts and decreases inversely proportional to the input power. Low frequency noise shows a $1/f$ dependence. © 2007 American Institute of Physics. [DOI: 10.1063/1.2722666]

Recent breakthroughs in cosmology have stimulated the development of a generation of detectors for millimeter and submillimeter wavelengths. The detectors should be fast and sensitive to perform a fast survey of the sky for measurements of, e.g., cosmic background radiation. The advanced cameras developed for astronomy usually find applications also in terahertz imaging for security, medical, and biological applications. For these applications, it is required to have fast and sensitive bolometers working at relatively high temperatures. A bolometer with a carbon nanotube (CNT) is an attractive option for these applications.

It was predicted in Ref. 1 that superconductor–insulator–normalmetal–insulator–superconductor (SINIS) structures can be effective bolometers and electron coolers of normal metal absorbers, similar to conventional semiconductor Peltier elements. Experimental cooling from 300 down to 100 mK in a SINIS structure was demonstrated in Ref. 2 and bolometric response up to 2 THz in Ref. 3. The characteristics of such devices can be estimated using the relation for energy transfer between electrons and phonons, $P = \Sigma \nu (T_e^5 - T_o^5)$, or for electron temperature, $T_e = (T_o^5 + P/\Sigma \nu)^{1/5}$, in which Σ is a material parameter, ν is the absorber volume, and T_o is the phonon temperature. According to these expressions, in order to increase the efficiency of warming by radiation and improve electron cooling, it is necessary to considerably reduce the volume of the absorber. In a conventional bolometer with a thin-film absorber fabricated using e-beam lithography,³ the volume is typically $10 \times 0.2 \times 0.05 = 0.1 \mu\text{m}^3$. It should be possible to reduce the volume of the absorber by replacing the metal with a carbon nanotube. In this way the volume could be reduced to $10 \times (0.0005)^2 \times 3.14 \approx 10^{-5} \mu\text{m}^3$, that is, four orders of magnitude smaller for an absorber of the same length.

Recently, a bolometric response to near-infrared radiation was reported for suspended macroscopic films of en-

tangled single-walled nanotubes.⁴ In that work, the nanotube samples had dimensions of $3.5 \times 0.5 \text{ mm}^2$ with a thickness in the range of 40 nm–1 μm . No bolometric response could be detected for nonsuspended nanotube films. In this letter, we report the results of a proof-of-principle experiment, where we show a strong bolometric response of deposited nanotube bundles to microwave radiation. This provides encouragement for the development of sensitive cold electron bolometers based on individual carbon nanotubes as the absorber, where one could expect considerably enhanced electron cooling to temperatures below 100 mK.

Our samples were fabricated on $7 \times 7 \text{ mm}^2$ oxidized *n*-doped silicon wafers. Ti/Au electrodes were first patterned using photolithography. An ac voltage is applied between the electrodes for dielectrophoretic deposition of single-walled nanotubes dispersed in H₂O with 4 wt % of sodium dodecyl sulphate. A relatively high nanotube concentration was deliberately used to deposit bundles of thousands of nanotubes rather than individual nanotubes between the electrode pairs. The voltage applied was 20 V peak to peak at a frequency of 13 MHz. After the nanotube deposition, four aluminum electrodes were deposited on top of the nanotube bundles. A typical device is shown in Fig. 1. The deposited nanotubes are expected to have a relatively high content of metallic nanotubes,⁵ but semiconducting nanotubes will also be present in the deposited bundles. Arrays of deposited nanotubes were preferred for these initial experiments instead of individual nanotubes, in order to reduce the resistance of the sample and enhance the response. The resistance of such samples at a temperature of 4 K was down to a few kiloOhms. The main part of this resistance seems to be due to the contact resistance.

The voltage response of such a CNT array to microwave radiation (110 GHz) was measured. The source of radiation was an impact ionization avalanche transit-time diode with an output power of about 10 mW. The signal was modulated with a mechanical chopper and attenuated with an EC-COSORB absorber. The radiation illuminated the CNT

^{a)}Electronic mail: mikhail.tarasov@mc2.chalmers.se

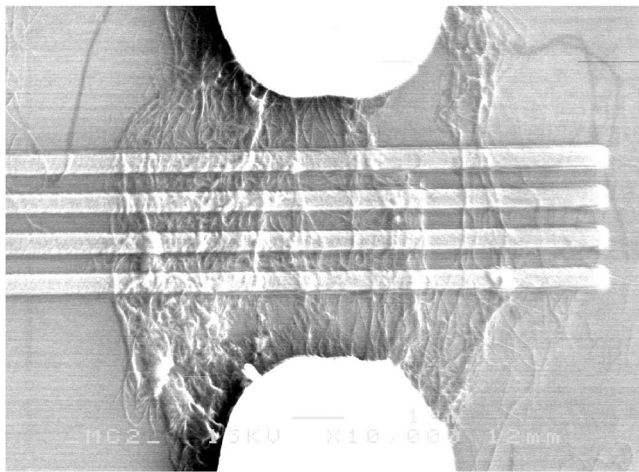


FIG. 1. Carbon nanotube bundles deposited across the top and bottom electrodes using ac dielectrophoresis. Four Al electrodes were patterned on top of the deposited nanotubes.

sample through a quasioptical window and cold IR filters. The sample was mounted on the flat side of an silicon extended hyperhemisphere silicon lens, 12.6 mm in diameter. The back side of the Si substrate was turned toward the silicon lens. The two of four outer Al electrodes (see Fig. 1) were dc biased and the voltage response on the two inner electrodes was lock-in detected, while the electrodes used for dielectrophoresis were kept unconnected. The I - V characteristics of the sample at 4.2 K, in the absence of microwave radiation, are shown in Fig. 2 (left hand axis) for a temperature of 4.2 K. The microwave response for a dc bias from -100 to $+100$ nA is shown on the same plot (right hand axis), demonstrating an output signal of over $50 \mu\text{V}$. The behavior looks very similar to the I - V curve and the response of a conventional SINIS bolometer with an absorber made of a normal metal thin-film strip.³ The difference is that in a SINIS bolometer, the nonlinearity originates from the superconductor energy gap of the SIN junction, and in the CNT junction, the nonlinearity is due to other effects such as Schottky-type junctions or Coulomb blockade.

The temperature dependence of the voltage on the same sample was measured at a bias current of $0.3 \mu\text{A}$ (see Fig. 3). From this dependence we can calculate a temperature response $dV/dT=0.4 \text{ mV/K}$ for $T < 50 \text{ K}$ that is close to the

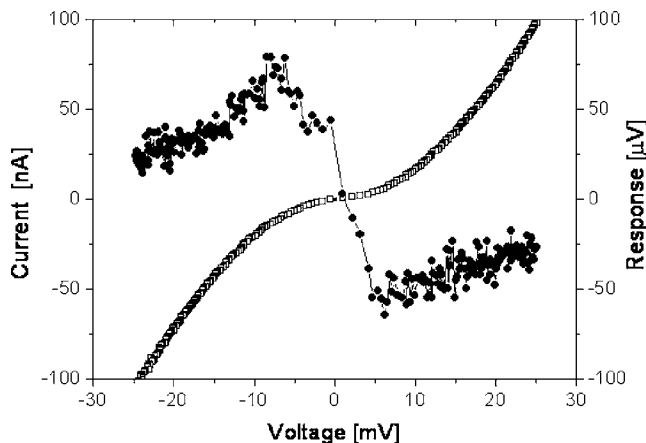


FIG. 2. I - V curve (left hand axis) measured between inner electrodes and voltage response of the nanotube device shown in Fig. 1, exposed to 110 GHz radiation. Measurements were performed at 4.2 K.

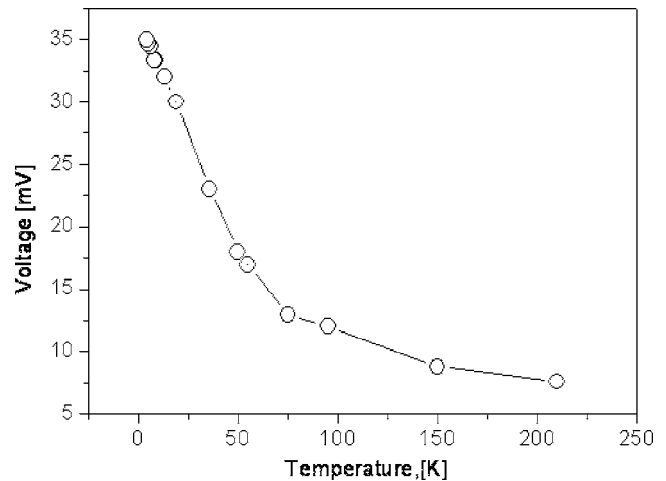


FIG. 3. Temperature dependence of the response voltage for a bias current of $0.3 \mu\text{A}$.

temperature response of the metallic thin-film SINIS bolometer in Ref. 3 (1.7 mV/K).

Due to a huge impedance mismatch (100Ω – $300 \text{ k}\Omega$) and high losses in the quasioptical beamguide, the incident power in the microwave measurements that is absorbed in the whole bundle can be very roughly estimated to be below 1 nW . To improve the accuracy of the responsivity measurements, we performed much more accurate ac measurements. The ac response was measured by applying an ac signal to the outer pair of the four electrodes and using the inner electrode pair as the detector at varying dc bias. The results are shown in Fig. 4 for an input power of 2.4 fW . There is a strong dependence of the response on the input signal power (Fig. 5). One can see that already for powers greater than a few femtowatts, the response decreases inversely proportional to the applied power.

To estimate the sensitivity, the noise of the device was also measured. It shows clear $1/f$ dependence (Fig. 6). From these measurements it is possible to deduce the noise equivalent power (NEP) by dividing the noise by the response. The best $\text{NEP}=3 \times 10^{-16} \text{ W/Hz}^{1/2}$ was measured at an input amplitude of 2.4 fW and an input frequency of 700 Hz , at

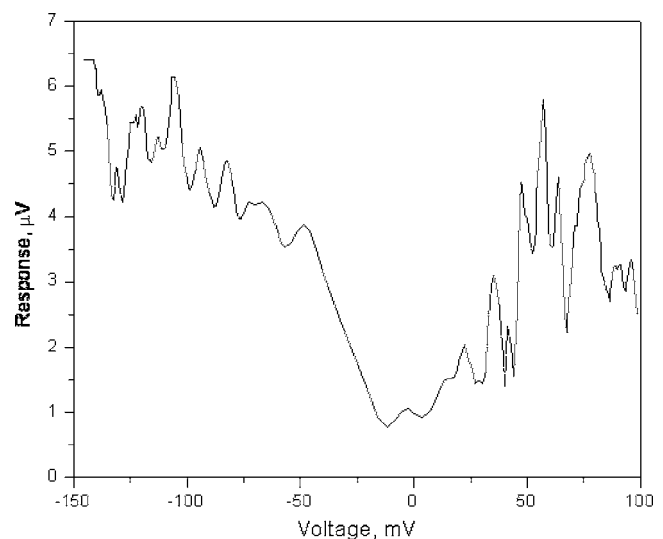


FIG. 4. ac response for input probe power of 2.4 fW . ac voltage applied across outer electrode pair; response measured across inner electrode pair.

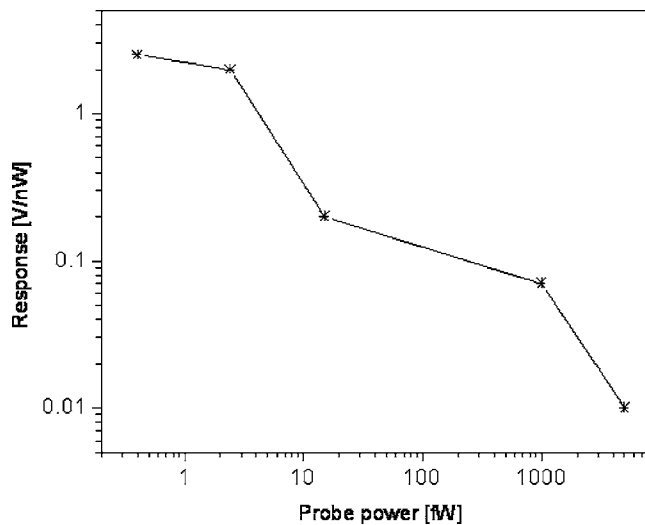


FIG. 5. Dependence of response on input probe power.

which the noise impact from the readout amplifier was the lowest.

We have demonstrated that carbon nanotubes can operate as bolometric detectors at microwave frequencies. Due to the small diameter of carbon nanotubes, the junction capaci-

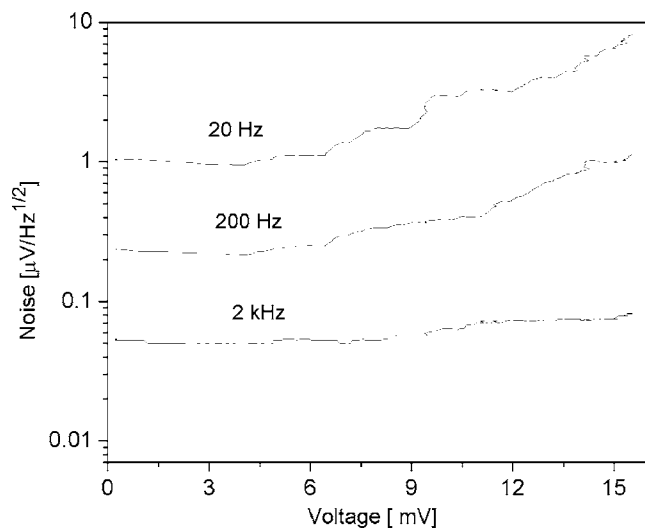


FIG. 6. Frequency dependence of noise in semilogarithmic plot.

tance can be significantly smaller than for normal Schottky detectors, which could push the detection frequency cutoff to beyond 30 THz.^{6,7} Besides bolometric detectors, other microwave applications of CNTs such as mixers,⁸ microwave field-effect transistors,⁹ and direct and heterodyne detectors¹⁰ have been realized. The NEP obtained in our work is better than what has been achieved with other direct detector device at an operating temperature of 4.2 K (see Ref. 11). However, it is important to emphasize that the measured NEP is for intrinsic electrical responsivity, while similar optical characteristics are many orders of magnitude worse. The fabrication technology should be further developed to improve impedance matching and quality of CNTs. The small volume of CNTs also makes them promising materials for the construction of cold electron bolometers formed from superconductor–insulator–normal metal structures.¹ Electron refrigeration from 360 down to 120 mK was recently demonstrated in a superconductor–insulator–normal metal structure,² and we expect to be able to improve this by using single CNTs or CNT bundles. However, for this purpose, the fabrication technology of the nanotube contacts will have to be improved.

This work was supported by the Swedish Institute, Vetenskapsrådet, K & A Wallenberg Foundation, and Grant Nos. PGLSS 7812.2006.2, and RFBR 05-02-19650.

- ¹L. Kuzmin, I. Devyatov, and D. Golubev, Proc. SPIE **3465**, 193 (1998).
- ²L. Kuzmin, I. Agulo, M. Fominsky, A. Savin, and M. Tarasov, Supercond. Sci. Technol. **17**, 400 (2004).
- ³M. Tarasov, L. Kuzmin, E. Stepanov, I. Agulo, A. Kalabukhov, Z. Ivanov, and T. Claeson, JETP Lett. **79**, 298 (2004).
- ⁴M. Itkis, F. Borondics, A. Yu, and R. Haddon, Science **312**, 413 (2006).
- ⁵D. S. Lee, D. W. Kim, H. S. Kim, S. W. Lee, S. H. Jhang, Y. W. Park, and E. E. B. Campbell, Appl. Phys. A: Mater. Sci. Process. **80**, 5 (2005).
- ⁶D. Tsang and S. Schwartz, Appl. Phys. Lett. **30**, 263 (1977).
- ⁷H. M. Manohara, E. W. Wong, E. Schlecht, B. D. Hunt, and P. H. Siegel, Nano Lett. **5**, 1469 (2005).
- ⁸S. Rosenblatt, H. Lin, V. Sazonova, S. Tiwari, and P. L. McEuen, Appl. Phys. Lett. **87**, 153111 (2005).
- ⁹A. A. Pesetski, J. E. Baumgardner, E. Folk, J. Przybysz, J. D. Adam, and H. Zhang, Appl. Phys. Lett. **88**, 113103 (2006).
- ¹⁰F. Rodriguez-Morales, R. Zannoni, J. Nicholson, M. Fischetti, K. S. Yngvesson, and J. Appenzeller, Appl. Phys. Lett. **89**, 083502 (2006).
- ¹¹A. Luukanen, E. N. Grossman, A. J. Miller, P. Heliö, J. S. Penttilä, H. Sipola, and H. Seppä, IEEE Microw. Wirel. Compon. Lett. **16**, 464 (2006).